Limits of declustering methods for disentangling exogenous from endogenous events in time series with foreshocks, main shocks, and aftershocks

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Many time series in natural and social sciences can be seen as resulting from an interplay between exogenous influences and an endogenous organization. We use a simple epidemic-type aftershock model of events occurring sequentially, in which future events are influenced (partially triggered) by past events to ask the question of how well can one disentangle the exogenous events from the endogenous ones. We apply both model-dependent and model-independent stochastic declustering methods to reconstruct the tree of ancestry and estimate key parameters. In contrast with previously reported positive results, we have to conclude that declustered catalogs are rather unreliable for the synthetic catalogs that we have investigated, which contains of the order of thousands of events, typical of realistic applications. The estimated rates of exogenous events suffer from large errors. The branching ratio *n*, quantifying the fraction of events that have been triggered by previous events, is also badly estimated in general from declustered catalogs. We find, however, that the errors tend to be smaller and perhaps acceptable in some cases for small triggering efficiency and branching ratios. The high level of randomness together with the long memory makes the stochastic reconstruction of trees of ancestry and the estimation of the key parameters perhaps intrinsically unreliable for long-memory processes. For shorter memories (larger "bare" Omori exponent), the results improve significantly.

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I. INTRODUCTION

A large variety of natural and social systems are characterized by a stochastic intermittent flow of sudden events: landslides, earthquakes, storms, floods, volcanic eruptions, biological extinctions, traffic gridlocks, power blackouts, breaking news, commercial blockbusters, financial crashes, economic crises, terrorist acts, geopolitical events, and so on. Sequences of such sudden events constitute often the most crucial features of the evolutionary dynamics of complex systems, both in terms of their description, characterization and understanding. Accordingly, a useful class of models of complex systems views their dynamics as a sequence of intermittent discrete short-lived events. In the limit where the time scales, over which the change of regimes associated with the occurrence of the events occur, are small compared with the intervent intervals, the catalog of events can be modeled using the mathematics of point processes [1,2]. This modeling strategy emphasizes that the system is active during short-lived events and inactive otherwise. This amounts to separating a more or less incoherent background activity (such as small undetectable earthquakes) from the occurrence of structured events (large earthquakes), which are the focus of interest. We note that the class of stochastic point processes is fundamentally different from that of discrete and continuous stochastic processes, for which the activity is nonzero most of the time.

It turns out that a significant understanding of the complex flow of observed events can be achieved by precisely framing the problem in terms of a classification of two limited classes of events: (i) those that are the response of the system to exogenous shocks to the system and (ii) those that

Having a time series or catalog of discrete events, we are interested in understanding the generating process that led to the observed sequence. The difficulty in deciphering the underlying mechanisms stems from the fact that the above systems of interest are on the one hand subjected to external forcing, which on the other hand provide them the stimuli to self-organize via negative and positive feedback mechanisms. Most natural and social systems are indeed continuously subjected to external stimulations, noises, shocks, solicitations, and forcing, which can widely vary in amplitude. It is thus not clear a priori if the observed activity is due to a strong exogenous shock, to the internal dynamics of the system organizing in response to the continuous flow of information and perturbations, or maybe to a combination of both. In general, a combination of external inputs and internal organization is at work and it seems hopeless to disentangle the different contributions to the observed collective human response. Determining the chain of causality for such questions requires disentangling interwoven exogenous and endogenous contributions with either no clear or too many signatures. How can one assert with confidence that a given event or characteristic is really due to an endogenous selforganization of the system, rather than to the response to an external shock?

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appear endogenously without any obvious external causes. This can be done by looking at the specific endogenous and exogenous signatures and their mutual relations, which are reminiscent of the fluctuation-susceptibility theorem in statistical physics [3,4]. This approach provides a useful framework for understanding many complex systems and has been successfully applied in several contexts: commercial book successes [5,6], social crises [7], financial volatility [8], financial bubbles and crashes [9,10], earthquakes [11,12], diseases in complex biological organisms [13], epileptic seizures [14], and so on.

A common feature observed in these different systems is the fact that events are not independent as they would be if generated by a Poisson process. Instead, they exhibit pronounced interdependencies, characterized by "selfexcitation," i.e., past events are found to often promote or trigger (in part) future events, leading to epidemiclike cascades of events. The analogy with triggering and cascade processes occurring in viral epidemics is so vivid in some instances that the name "epidemic-type aftershock (ETAS) model" has been given to one of the most popular model of earthquake aftershock processes [15–17]. The ETAS model belongs to the class of self-excited conditional point processes introduced in mathematics by Hawkes [18-20]. It constitutes an excellent first-order approximation to describe the spatiotemporal organization of earthquakes [21] and is now taken as a standard benchmark. This class of "selfexcited" point processes also provides quantitative predictions on the different decay rates after exogenous peaks of activity on the one hand and endogenous peaks of activity on the other hand. These predictions have been verified in a unique data set of almost 5 000 000 time series of human activity collected subdaily over 18 months since April 2007 from the third most visited website YouTube.com.

Given these preliminary successes, we would like to decipher, understand, and perhaps forecast the dynamics of events. For this, it is important to recognize that the observed dynamics can be modeled as a complex entangled mixture of events, from exogenous shocks impacting the systems providing bring surprises, that are progressively endogenized by the system, which is also capable of purely endogenous (or internal) bursts of activity. In between, real systems can be viewed as organized by a mixture of exogenous shocks and endogenous bursts of activity. The grail is to disentangle these different types of events.

The purpose of the present paper is to contribute a step toward the operational problem of disentangling the exogenous and endogenous contributions to the organization of a system revealed by a time series of discrete events. For this, we use synthetic catalogs of events generated by the ETAS model. This model describes the occurrence of successive events, each of them characterized by its time t_i of occurrence and a "mark" M_i , which we refer to as "magnitude" to borrow from the vocabulary of earthquakes. Generally, the mark can be any trait or property that influences the ability of the event to trigger other future events. The conditional intensity $\lambda(t, M | H_t)$ of the linear version of the general self-excited Hawkes process (which we use later to define the specifics of the ETAS model) reads

$$\lambda(t, M | H_t) = \mu(t, M) + \sum_{t_i < t} h(t, M; t_i, M_i).$$
(1)

 $\lambda(t, M | H_t)$ consists of two main contributions: (i) the background intensity μ which, if alone, would give a pure Poisson process of background events (more complicated source terms can be considered, but we keep μ constant in the present paper and thus drop its possible dependence on t and M); (ii) the response functions h, one for each past event, describing the propensity to trigger future events. Specifically, $h(t,M;t_i,M_i)$ is the intensity of the *i*th event that occurred at time t_i with magnitude M_i to produce an aftershock at a later time $t > t_i$ with magnitude M. H_t stands for the history known at time t, and thus includes the sequence of all events from the beginning of observations till t.

Zhuang et al. [22,23] proposed a rigorous stochastic declustering method that, in essence, implements the program of disentangling the different events according to their background (exogenous) or triggered (endogenous) origins. Zhuang et al. applied the space-time ETAS model to real earthquake catalogs. The main problem is that no validation test was performed to check if the reconstruction of the ancestry sequence was realistic and if the inverted parameters were consistent, i.e., without bias. In the present paper, our goal is to make a systematic investigation of the quality of the stochastic reconstruction method(s). For this, using ETAS model (1), we generate known synthetic catalogs of events, which are considered as our "observations." We then apply different approaches, which are variants of the "stochastic declustering" methods introduced by Zhuang et al. [22,24] and generalized by Marsan and Lengliné [25]. Our strategy is to compare the key parameters obtained from the reconstruction of the cascade of triggering events obtained by these declustering methods applied to our "observations" to the known true values. This allows us to establish the uncertainties and biases associated with these "model inversions." As we shall see, the level of stochasticity, inherent in self-excited conditional Poisson processes, introduces rather dramatic errors, which appear to have been largely underestimated in the literature. Because of the general application of this problem to a variety of domains, we focus our attention to time series of events, assuming that spatial information is either irrelevant or not available. This makes the problem of stochastic declustering less constrained than in the case of earthquakes, for instance, in which one does have some information on the spatial positions, in addition to the times of occurrence.

The paper is organized as follows. Section II reviews the two methods of stochastic declustering. Section III first describes the ETAS model that we have used for the generation of synthetic catalogs. Then, it presents the two specific implementations of the ETAS model, referred to respectively as the generation-by-generation and event-by-event algorithms. After recalling the main results of previous tests performed by other authors, Sec. IV defines the parameters that are tested and presents the main results using the two declustering stochastic declustering method (SDM) and Marsan and Lengliné's model-independent stochastic declustering (MISD) methods on synthetic catalogs generated by the two generation-by-generation and event-by-event algorithms. Section VI concludes.

II. DESCRIPTION OF DIFFERENT DECLUSTERING VARIANTS

The general idea underlying a stochastic declustering method applied to a sequence of events thought to be generated by a self-excited conditional Poisson process is to attribute to each event a probability of being either an exogenous (a so-called "background" event) or an offspring of previous events (endogenous). Obviously the former probability is one minus the later probability. Therefore the main task is to estimate one of these probabilities for each event. Specifically, it is convenient to focus on the probability that a given event is a background event.

We start from the stochastic declustering method introduced in Refs. [22,24]. This method can be implemented under several technical variants, such as the thinning (random deletion) method, and a variable bandwidth kernel function method associated with the maximum-likelihood estimation of the ETAS model. Using any of these two approaches, the background intensity is obtained by using an iterative algorithm, and the declustered catalogs can also be generated.

A. Zhuang et al.'s declustering method

We focus on the thinning procedure, which uses the probabilities ρ_{ik} for the *k*th event to be an aftershock of the *i*th event and the probability ϕ_k that the *k*th event is only a background event [22,24]. These probabilities can be expressed in terms of the response function and the intensity defining ETAS model (1),

$$\rho_{ik} = \frac{h(t_k, M_k; t_i, M_i)}{\lambda(t_k, M_k | H_t)}.$$
(2)

This allows us to define the probability ρ_k that the *k*th event is an aftershock (whatever its triggering "mother") as

$$\rho_k = \sum \rho_{ik}.$$
 (3)

Therefore, the probability ϕ_k that the *k*th event is only a background event is

$$\phi_{k} = 1 - \rho_{k} = \frac{\mu(t_{k}, M_{k} | H_{t_{k}})}{\lambda(t_{k}, M_{k} | H_{t})}, \qquad (4)$$

These probabilities are called "thinning probabilities" [23].

If we delete the *k*th event in the catalog with probability ρ_k for all $k=1;2;\ldots;N$, then the thinned process should realize a nonhomogeneous Poisson process of the intensity $\mu(t,M)$ (see [26] for the mathematical justification). This process is called the background subprocess, and the complementary subprocess is the cluster or offspring process. The following algorithm implements this thinning procedure.

Algorithm 1

The indices k=1;2;...;N of the events in the catalog are ordered according to their time sequence: $t_i < t_2 < ... < t_N$.

- For all events k=1,2,...,N;, calculate their probabilities ρ_k in Eq. (3).
- (2) Generate N uniform random numbers U_1, U_2, \ldots, U_N in [0;1].
- (3) If U_k < 1-ρ_k, keep the kth event; otherwise, delete it from the catalog as an offspring. The remaining events can be regarded as the background events.

This algorithm can be applied to any data series and will find thinning probabilities for each event, which are functions of the specific model used.

The next issue is to estimate the parameters defining the response functions $h(t_k, M_k; t_i, M_i)$ of the Hawkes point process. For this, the following algorithm determine which event in the data set is the ancestor of a given *k*th event.

Algorithm 2

- For each pair of events *i*;*k*=1,2,...,*N*(*i*<*k*), calculate the probability ρ_{ik} in Eq. (2) and φ_k in Eq. (4).
- (2) Set k=1.
- (3) Generate a uniform random number $U_k \in [0; 1]$.
- (4) If U_k < φ_k, then the kth event is regarded as a background event.
- (5) Otherwise, select the smallest index *I* in {*i*+1,...,*N*} such that U_k < φ_k+Σ^I_{i=1} ρ_{ik}. Then, the *k*th event is regarded to be a descendant of the *I*th event.
- (6) If k=N, terminate the algorithm; else set k=k+1 and go to step 3.

The output of this algorithm is to provide a complete classification of events as backgrounds or descendent of some previous event, this previous event being either a background or a descendent of another previous event and so on. The corresponding reconstructed ancestry tree allows us to calculate different parameters of the model, such at the productivity law, giving the average number of offsprings of a given event as a function of its magnitude. Of course, a given ancestry tree obtained by this stochastic declustering method is not unique since, given a fixed catalog of events, it depends on the realization of the random numbers $\{U_k\}$ in Algorithms 1 and 2. Thus, these stochastic reconstructions must be performed many times with different independent realizations of the random numbers $\{U_k\}$ to obtain many statistically equivalent ancestry trees over which statistical averages can be taken.

B. Marsan and Lengliné's model-independent stochastic declustering

The MISD method also aims at determining thinning probabilities (2) [25,27], using a rapidly converging algorithm with a minimum set of hypotheses. The MISD method proposes to reveal the full branching structure of the triggering process, while avoiding model-dependent inversions.

The MISD method has the following key assumptions:

(i) The generating process of the observed catalog of events is considered to be a point process, in time, space, and magnitude (we will only consider the situation where no spatial information is provided in this paper to focus only on time series). Its conditional intensity consists of a linear superposition of a constant (Poisson process) background intensity μ and of the branching part represented by the summation in the right-hand side of the following expression:

$$\lambda(t, \overline{x}, m) = \mu + \sum_{t_i \le t} \lambda_i(t_i, \overline{x}_i, m_i), \qquad (5)$$

where the sum is performed over all ith events that have occurred before time t.

(ii) The average activity of events in response to the occurrence of a given event only depends on its magnitude m.

While very general, it is however necessary to point out that the MISD method assumes that the conditional Poisson intensity $\lambda_i(t_i, \bar{x}_i, m_i)$ due to each passed event *i* contributes additively (linearly) to the total conditional intensity $\lambda(t, \bar{x})$. This assumption is appropriate if the generating process belongs to the class of linear self-excited Hawkes processes (1). However, this linearity condition excludes a large class of nonlinear self-excited conditional Poisson models [28,29], and in particular the class of multiplicative (in contrast to additive) Poisson models [30,31] endowed with general multifractal scaling properties [32].

The MISD algorithm includes two iterating steps (that we present in full generality, including the possible existence of a spatial information):

(1) Starting with a first *a priori* guess of the "bare" kernel $\lambda(t, \bar{x}, m)$ and of μ , the triggering weight (thinning probability) is estimated by $\rho_{ij}=a_j\lambda(t_j-t_i, |\bar{x}_j-\bar{x}_i|, m_i)$ for $t_i < t_j$ and 0 otherwise. The "background weight" is estimated as $\phi_i=a_i\mu$, where a_i is a normalization coefficient.

(2) The updated (*a posteriori*) "bare" rates are then computed as

$$\lambda(|\Delta t|, |\Delta \overline{x}|, m) = \frac{1}{N_m \delta t \, S(|\Delta \overline{x}| \, \delta r)} \sum_{i, j \in A} \rho_{ij},$$

where A is the set of pairs such that $|\bar{x}_i - \bar{x}_j| = |\Delta \bar{x}| \pm \delta r$, $m_i = m \pm \delta m$ and $t_j - t_i = t \pm \delta t$ (δr , δt , and δm are discretization parameters), N_m is the number of events such that $m_i = m \pm \delta m$, and $S(|\Delta \bar{x}| \delta r)$ is the surface covered by the disk with radii $|\Delta \bar{x}| \pm \delta r$. Similarly, the *a posteriori* background rate is

$$\mu = \frac{1}{TS} \sum_{j=1}^{N} \phi_j,$$

where T is the duration of the time series (containing N events) and S is the surface analyzed.

Roughly speaking, the first step of the algorithm selects the triggering events for each triggered event (i.e., it assigns triggering weights based on our present knowledge of the rates). The second step then updates these rates, using the intermediate branching structure obtained at the first step. The solution is accepted (convergence is achieved) when the *a priori* and the *a posteriori* kernels are identical, implying that the rates and weights are consistent with each other.

The initial formulation recalled above and the implementation of the MISD algorithm was done for three- and fourdimensional data, including time, magnitude, and spatial coordinates of events [25,27]. Lately, Marsan adapted the code to two-dimensional data (times and magnitudes) and we use his code for the research reported here. In our implementation, following Marsan and others, we use logarithmic binned time and linear magnitude bins.

III. MODEL AND SIMULATION IMPLEMENTATIONS

This section describes the specific Hawkes-ETAS model that we use to generate synthetic time series of events, which are then collated in catalogs on which the different stochastic declustering algorithms described in the previous section are applied. It should be noted that the Hawkes-ETAS model includes many point processes as particular cases. It constitutes a very general class of "linear" point processes. The term "linear" refers here to the dependence of intensity (1) as a (linear) sum over the past events.

Previous implementation by Zhuang *et al.* [22,24] have used the ETAS model with its full space-time formulation. Here, we use the ETAS model in which the spatial information is supposed to be nonexistent or irrelevant. In this way, our tests are relevant to catalogs of events obtained in other systems, such as social, commercial financial, and biological systems.

A. ETAS model

The ETAS model is defined as the self-excited linear conditional Poisson model with intensity (1), in which the ("bare") response function $h(t, M; t_i, M_i)$ is expressed at the product of three terms,

$$h(t, M, t_i, M_i | H_{t_i}) = j(M)\Phi(t - t_i | H_{t_i})Q(M_i),$$
(6)

where

$$j(M) = b \left[\ln(10) \right] \, 10^{-b(M-M_0)} \tag{7}$$

has the form of a Gutenberg-Richter law prescribing the frequency of events of magnitudes M,

$$\Phi(\tau) = \frac{\theta c^{\theta}}{(c+\tau)^{1+\theta}},\tag{8}$$

has the form of the Omori-Utsu law specifying the probability density function (PDF) of time intervals between a main event and its direct aftershocks, and

$$Q(M) = K \ 10^{\alpha M} \tag{9}$$

is the productivity law giving the mean number of direct aftershocks generated by an event as a function of its magnitude M. We thus have

$$h(t, M, t_i, M_i | H_{t_k}) = b \left[\ln(10) \right] 10^{-b(M-M_0)} \cdot \frac{\theta c^{\theta}}{(c+t-t_i)^{1+\theta}} K 10^{\alpha M_i}.$$
(10)

Equations (6) and (10) express the independence between the determination of the magnitude of aftershocks and their occurrence times on the one hand, and with the magnitudes of their triggering ancestors on the other hand. Next, the background rate $\mu(t, M)$ in Eq. (1) is taken also multiplicative as

$$\mu(t,M) = j(M)\mu. \tag{11}$$

The constant μ means that the background events are occurring according to a standard memoryless Poisson process with constant intensity. The multiplicative structure of $\mu(t,M)$ in Eq. (11) again expresses that the magnitudes of the background events are independent of their occurrence times.

In summary, the conditional intensity of the ETAS model used here reads

$$\lambda(t, M | H_{t_k}) = j(M) \left(\mu + \sum_{i=1}^k \Phi(t_i | H_{t_k}) Q(M_i) \right), \quad (12)$$

with definitions (7)–(9) and the constant background rate μ . Substituting Eqs. (11) and (12) into Eqs. (3) and (4), we get the sought thinning probabilities for the ETAS model.

In order to generate synthetic catalogs of events, we considered two different simulation algorithms. Both techniques used the following set of parameters:

$$\Theta = (c, \theta, \alpha, b, n), \tag{13}$$

where

$$n = \frac{K}{1 - \alpha/b} \tag{14}$$

is the branching ratio—the mean number of direct aftershocks per triggering event [16]. The first algorithm keeps the information about the mother-daughter relations by generating events generation by generation. It uses a slightly different implementation of the ETAS model than the second algorithm and is computationally more costly in the sense that a generated catalog becomes stationary only after tens or even hundreds of thousands events. The second algorithm, which is based on the formulation of the model described in Sec. III A, is not very fast but is efficient. It generates events that belong to the stationary regime from the beginning, so no information about preceding events is lost (unlike the first algorithm). However, it loses the information about the mother-daughter relations. Its performance was validated in Ref. [33].

B. First algorithm (generation by generation): Generation of synthetic catalogs keeping the information on the relations between events

The first simulation algorithm we use here was developed by Felzer [34]. The idea is to generate events generation by generation. First, the mother-shock and the background events are generated, then goes the first generation of aftershocks of existing events, after that the second generation and so on. The procedure stops when the time boundary or the limit on the number of events is reached. We slightly modified the algorithm (mostly input and output) to correspond to our needs.

The advantage of knowing the ancestry relations between events, which allows us to estimate parameters such as α and b, comes at the cost of the existence of a long transient before the time series of events become stationary. Exactly at criticality n=1, where n is given by Eq. (14), this transient becomes infinitely long lived since the renormalized Omori law [16], which takes into account all generations of events triggered by each event, develops a nonintegrable decay $1/t^{1-\theta}$, where θ is defined in Eq. (8). This makes this first algorithm unreliable for n close to 1.

C. Second algorithm (event by event): Generation of synthetic catalogs without any information on the relations between events

Compared with standard numerical codes that have been used by previous workers to generate synthetic catalogs of events, the event-by-event algorithm uses the specific formulation of the ETAS model and of its known conditional cumulative distribution function (CDF) of interevent times. Thus, parts of the calculations can be done analytically. The occurrence times are generated one by one by determining the CDF of the time till the next event based on the knowledge of the previous CDF and the time of the previous event (first introduced by Ozaki for Hawkes' processes [35]). For this, we use a standard Newton algorithm as well as a standard randomization algorithm. The event-by-event algorithm is not very fast because it has to solve the equation for the CDF numerically each time. However, as already mentioned above, its advantage stems from the fact that the generated events belong to the stationary regime from the beginning of each catalog.

Let us recall how the CDFs $F_{k+1}(\tau)$ for k=2,3,... of successive events are obtained recursively. The first event is supposed to have occurred at time $t_1=0$.

The CDF $F_2(\tau)$ of the waiting time from the first to the second shock is made of two contributions: (i) the second shock may be a background event or (ii) it may be triggered by the first shock. This yields

$$F_2(\tau) = 1 - e^{-\omega\tau} e^{-q_1[1 - a(\tau)]},\tag{15}$$

where $q_1 = Q(m_1)$ is the productivity of the first shock obtained from expression (9) given its magnitude m_1 , and $a(\tau)$ is defined as

$$a(\tau) = \int_{\tau}^{\infty} \Phi(t') dt' = \left(1 + \frac{\tau}{c}\right)^{-\theta}.$$
 (16)

All following shocks are similarly either a background event or triggered by one of the preceding events. The CDF $F_3(\tau)$ of the waiting time between the second and the third shocks is thus given by

$$F_3(\tau) = 1 - e^{-\omega\tau} e^{-q_1 [a(t_2 - t_1) - a(t_2 + \tau - t_1)] - q_2 [1 - a(\tau)]}, \qquad (17)$$

where t_2 is the realized occurrence time of the second shock. Iterating, we obtain the CDF $F_k(\tau)$ for the waiting time between the (k-1)th and kth shocks under the following form:

$$F_{k}(\tau) = 1 - e^{-\omega\tau} \exp\left[-\sum_{i=1}^{k-1} q_{i}(a(t_{k-1} - t_{i}) - a(t_{k-1} + \tau - t_{i}))\right],$$
(18)

where t_j is the occurrence time of the *j*th event (which is equal to sum of all generated time intervals between events prior to the *j*th one), and $q_i = Q(m_i)$ is the productivity of the *i*th shock obtained from expression (9) given its magnitude m_i .

In order to generate the (k+1)th interevent time interval between the occurrence of the kth and (k+1)th shock, it is necessary to know the k previous interevents times between the k previous shocks and their k magnitudes. Since, in the ETAS model, the magnitudes are drawn independently according to Gutenberg-Richter distribution (7), they can be generated once for all. In order to generate a catalog of Nevents, we thus draw N magnitudes from law (7). In order to generate the corresponding N interevent times, we use expression (18) iteratively from k=1 to k=N in a standard way: since any CDF F(x) of a random variable x is by construction itself uniformly distributed in [0,1], we obtain a given realization x^* of the random variable x by drawing a random number r uniformly in [0,1] and by solving the equation $F(x^*) = r$. In our case, we generate N independent uniformly distributed random numbers x_1, \ldots, x_N in [0,1] and determine each τ_i successively as the solution of $F_i(\tau_i) = x_i$.

As mentioned above, the main shortcoming of this procedure is that it does not record if an event was spontaneous or a descendant of some previous event. For that reason, we cannot use it for all parameter estimations.

D. Preliminary tests of the synthetic catalogs

We have checked the consistency of our algorithms by verifying that, for K=0 in Eq. (9) corresponding to the absence of triggering, a Poisson flow of event time occurrence is obtained.

We have implemented the two algorithms just discussed in the previous subsection and have constructed the corresponding CDFs from the obtained time series for b=1, $c = 10^{-3}$, $\alpha=0.7$, $\theta=0.1$, n=0.7 (K=0.21), and $m_d-m_0=0.01$. Figure 1 shows that both algorithms lead to CDFs which are very close to each other, when the transient regimes of the catalogs generated by the first generation-by-generation algorithm are removed. Using the whole catalog including the transient part for the first algorithm leads to very large distortions. We interpret the remaining slight difference between the CDFs of interevent times for the first and second algorithms after (partial) removal of the transient as due to the residual influence of the transient part of the catalog in the first generation-by-generation algorithm IIIB.

Some words of caution are in order on the meaning of "transient." From a pure theoretical view point, it is actually



FIG. 1. CDFs of interevent times for synthetic catalogs generated for b=1, $c=10^{-3}$, $\alpha=0.7$, $\theta=0.1$, n=0.7 (K=0.21), and m_d $-m_0=0.01$. The CDFs of interevent times are shown for the two algorithms (first event-by-event IIIC and second generation-bygeneration IIIB) and for the Poisson flow of the background events, bare Omori law (8), and the theoretical prediction obtained from the linearized equation governing the generating probability function of the interevent CDFs of the ETAS model developed in [21,36] (referred to in the legend as "CDF found from linearized ETAS model").

intrinsically impossible to remove absolutely the transient regime when the exponent θ of bare response function (8) is smaller than 1. Indeed, the statistical expectation $\langle \tau \rangle$ of the waiting time between a main event and its direct aftershocks is infinite: $\int_0^T \Phi(\tau) \tau d\tau \sim T^{1-\theta}$ diverges as one considers larger and larger catalogs of increasing durations T to perform the average. This regime is typical of processes with "infinite" memory and is well known to lead to anomalous transport and diffusion properties (see for instance Refs. [37-40]). In practice, we have experimented with different definitions of the transient regime to find an operational procedure: typically, when we remove the first half of the catalogs, we found that the CDFs reconstructed from the data in the second half of the catalogs were close to the exact distributions and did not change appreciably by removing even more events. Having said that, the infinite memory associated with the range $0 < \theta \le 1$, often documented in empirical data, is one main cause of the difficulties with inverting for the parameters, as we demonstrate below. Specifically, Sec. V shows that the stochastic reconstruction methods work much better for $\theta > 1$.

Figure 1 also shows for comparison the interevent CDFs of (i) the Poisson flow of the background events, (ii) bare Omori law (8), and (iii) the theoretical prediction obtained from the linearized equation governing the generating probability function of the interevent CDFs of the ETAS model developed in [21,36], confirming the need for the full non-linear form of the equation of the generating probability function of the interevent CDFs [33].

TABLE I. Estimated parameters of the ETAS model obtained by the stochastic declustering method (SDM) of Zhuang *et al.* [22,24,41] described in Sec. II A applied to real earthquake data from New Zealand (NZ), Central and Western Japan (CJ and WJ), and Northern China (NC). The quoted values for the fertility exponent α differ from those reported by Zhuang *et al.* by the conversion factor ln 10 accounting for our use of base-ten logarithm and exponential compared with the natural logarithm and exponential used by Zhuang *et al.*

Region	С	θ	α	K	п
NZ	0.017	0.164	0.389	0.335	0.548
CJ, WJ	0.040	0.250	0.495	0.204	0.404
NC	0.003	0.030	0.499	0.546	1.090

IV. RESULTS AND PERFORMANCE OF STOCHASTIC DECLUSTERING METHODS (SDM AND MISD)

A. Previous tests of Zhuang et al.

As mentioned before, Zhuang *et al.* [22,24,41] applied their declustering procedure described in Sec. II A to real earthquake catalogs over four geographical regions: New Zealand (NZ), Central and Western Japan (CJ and WJ), and Northern China (NC). Table I provides the results of their SDM applied to these four regions.

Our main remark is that no error or uncertainty analysis is reported and no study of the impact of the lower magnitude threshold used in the catalog is performed. This is particularly worrisome, given the demonstration that parameter estimations are significantly biased when the minimal observable (registerable) magnitude m_d is different from (usually larger than) the minimal event magnitude m_0 able to produce aftershocks [42,43].

In a later paper, Zhuang *et al.* [44] reported some synthetic tests to assess the reliability of the SDM in the ideal case where $m_d = m_0$. They quoted "good reconstruction results." However, the parameters estimated with their SDM $(c \approx 0.0004, \ \alpha \approx 0.57, \ \theta \approx 0.014, \ n \approx 0.25)$ were very different from the true parameters $(c=0.0002, \ \alpha \approx 0.65, \ \theta = 0.12, \ n \approx 0.99)$ used to generate the synthetic catalogs. Very worrisome is the very large error in the value of the branching ratio n. The true value n=0.99 corresponds to a system close to critical branching in which triggering of multiple generations is expected to be very strong since 99% of events are triggered on average while only 1% are exogenous [12]. In contrast, the estimated value $n \approx 0.25$ would be interpreted as a relatively weak triggering regime in which three-quarters of the events are exogenous.

B. Previous tests of Marsan and Lengliné

Unlike the SDM of Zhuang *et al.*, the MISD method was partially verified. Marsan and Lengliné generated synthetic catalogs with the parameters b=1, c=0.01, $\alpha=0.87$, $\theta=0.2$, n=0.9, and $\mu=0.25$. Applying the MISD to those catalogs, Marsan and Lengliné could estimate the background rate $\mu^{est}=0.248\pm0.01$, very close to the true value. The estimates of the other parameters were reported to be also good [25].

TABLE II. Direct estimation of the parameters Θ^s with the modified SDM compared with the true parameters Θ^d of the ETAS model used to generate the synthetic catalogs by the eventby-event algorithm.

No.		b	α	С	θ	п
1	Θ^s	1.00	0.20	0.001 00	0.50	0.20
	Θ^d	3.80	0.27	0.001 03	0.54	0.07
2	Θ^s	1.00	0.50	0.001 00	0.50	0.50
	Θ^d	0.44	0.30	0.000 50	0.21	0.78
3	Θ^s	1.00	0.80	0.001 00	0.50	0.80
	Θ^d	1.00	0.96	0.001 60	0.64	0.96
4	Θ^s	1.00	0.20	0.001 00	0.50	0.80
	Θ^d	0.33	0.24	0.002 40	0.84	0.92
5	Θ^s	1.00	0.80	0.001 00	0.50	0.20
	Θ^d	1.00	0.76	0.000 70	0.41	0.10

C. Self-consistency of parameter estimations of n and α from synthetic catalogs

1. Modifications to SDM

In the above presentation, we have considered only the thinning probabilities of events in a catalog. But Zhuang *et al.*'s method contains in addition a determination of conditional intensity function (12) obtained by using an iterative search procedure with the maximum likelihood $L(\Theta)$ [2]. Specifically, the search procedure determines the set of parameters Θ of the model for which the thinning probabilities are best consistent with the conditional intensity $\lambda(t, M | H_{t_k})$. Zhuang *et al.* used the first-order algorithm of Davidson-Fletcher-Powell to find the sought maximum of the likelihood function. This method suffers from the need to calculate the derivatives of $\lambda(t, M | H_{t_k})$ and $L(\Theta)$, which makes errors accumulate and slows down the calculations.

In the present work, we use the Nelder-Mead simplex method which is significantly more efficient than the firstorder Davidson-Fletcher-Powell algorithm. In addition, we recalculate the probabilities ρ_{ij} for each estimation of the likelihood function rather than at each iteration as performed in the SDM version used by Zhuang *et al.* Our approach increases the computation time of a single iteration but actually provides a significant gain as the number of iterations needed for convergence is greatly reduced.

The set of parameters $\Theta^{\overline{d}}$ obtained from the maximization of the likelihood function corresponds to the *direct estimation*. The priority is to verify whether the direct estimation is good enough. Then, we need to check that the catalog reconstructed using thinning probabilities also provides good estimates of the model parameters.

2. Applying the modified SDM to two-dimensional data

To check the goodness of the direct estimation by the modified SDM (mSDM), we generated five catalogs of 2500 events for different values of α and n and fixed c=0.001, b=1, and $\theta=0.5$. As one can see from the results presented in Table II, the direct estimation can be quite far off, in particu-

lar for the parameters θ and *n*. But the errors turn out to be smaller than with the thinning probabilities that will be presented below. One of the origins for the errors is the limited length of the synthetic catalog, notwithstanding the use of an intentionally large value of $\theta = 1/2$ leading to a rather short memory (compared to that for smaller values of θ used below). The SDM and mSDM need more events to reduce the errors in the parameter estimation. Further below, we will consider the relationship between the length of a catalog that is needed to obtain reasonable results and the memory quantified by the exponent θ).

3. Targeted parameters

As a first test, we assume known the parameters of the ETAS process generating the synthetic catalogs. For instance, we can assume that the use of the mSDM led us to the true values of parameters $\Theta^d \approx \Theta^s$. We then apply the SDM and the MISD to determine the background and triggered events. From this knowledge, we can estimate directly the branching ratio *n* and the fertility exponent α and compare them with the true values. We stress that we use the exact parameters that enter in the generation of the synthetic catalogs to find thinning probabilities (2) to perform this test. Thus, any discrepancy between the estimation of *n* and α using the thinning probabilities should be attributed only to errors in the reconstruction of the tree of ancestry through the thinning probabilities.

The first key parameter of interest is the branching ratio, which can be estimated from the knowledge of the number of exogenous events within the data set, according to

$$n_e = 1 - \frac{\text{number of background events}}{\text{total number of events}},$$
 (19)

where the subscript e indicates that n_e is experimental or estimated value.

Knowing the tree structure of events obtained from the declustering method, we can estimate directly productivity law (9). Specifically, the tree structure allows us to calculate straightforwardly the mean number of direct aftershocks triggered by a given event, and then to test how this mean number depends on the magnitude of the mother event. Given true law (9), the estimated dependence of the number of aftershocks as a function of the magnitude of the main shock is fitted by the following expression:

$$\tilde{Q}(M) = K \ 10^{A^*M},$$
 (20)

where (K^*, A^*) are determined using standard optimization algorithms. The estimated values can then be compared to the true values (K, α) used to generate the synthetic catalogs.

To account for the fact that, in real time series, small events below a magnitude detection threshold $m_d > m_0$ are not detected, we also investigate the influence of this detection threshold on the estimated parameters. Intuitively, as demonstrated in Refs. [42,43], missing events lead to misinterpret triggered events as exogenous background events since the chain of causal triggering may be ruptured. This may influence severely the estimation of the background rate and therefore of the parameters controlling the fertility and triggering efficiency of past events. A good diagnostic of the effect of missed events is the branching ratio n [42,43]. We thus evaluate the apparent branching ratio estimated by the declustering method on the catalog of events with magnitudes larger than m_d according to the following formula:

$$n_t = \frac{\text{number of aftershocks with } M > m_d, \text{ whose mother also has } M > m_d}{\text{number of all events with } M > m_d}$$
(21)

We also will test how the truncation affects the estimated parameters, and does this correspond to the theoretical dependence [21,36].

4. First test on declustering Poisson sequences

We first applied the SDM and MISD to catalogs generated by a simple Poisson process, obtained from the formulation of Sec. IV A by imposing the value K=0 in Eq. (9). As a consequence, only exogenous background events without any triggered event are generated in synthetic catalogs with intensity imposed equal to $\mu=1$. A correct declustering algorithm should find n=0 and a background rate $\hat{\mu}=1$.

Applying the SDM on tens of catalogs each containing between 1000 and 2000 events with magnitudes M > 0, we recovered the correct result that the branching ratio is n=0 in every case (all probabilities ϕ_i are found exactly equal to 1)

and the background rate is estimated as $\hat{\mu}=1\pm0.05$, close and consistent with the true value.

On a technical note, the arrest criterion controlling the convergence of the algorithm was found to play a strong role, much more significant for time-domain-only catalogs than when the catalogs include spatial information. We found these good results only for an arrest criterion corresponding to differences smaller than 0.0001 between the parameter values of successive iterations. We kept this value for all subsequent tests, as a compromise between accuracy of the convergence and numerical feasibility.

On the same catalogs, the MISD method found, however, nonzero probabilities ρ_{ij} for $i \neq j$ and, as a result, a nonzero and actually quite large branching ratio $n=0.23\pm0.03$. The estimation of the background rate was also not very good $\hat{\mu}_0=0.81\pm0.02$ instead of the true value 1. This means that the MISD method applied in the temporal domain to already

argonum. The other parameters have been fixed to $m_d = m_0 = 0$, $b = 1$, $c = 0.002$, and $b = 0.19$.								
No.	п	n _e	α	A^*	K	Κ*		
1	0.26	0.237 ± 0.009	0.0	0.024 ± 0.041	0.260	0.930 ± 0.097		
2	0.26	0.239 ± 0.010	0.1	0.015 ± 0.060	0.234	0.939 ± 0.118		
3	0.26	0.240 ± 0.010	0.2	-0.005 ± 0.041	0.208	0.960 ± 0.087		
4	0.26	0.237 ± 0.009	0.3	-0.048 ± 0.039	0.182	1.034 ± 0.079		
5	0.26	0.240 ± 0.012	0.4	-0.059 ± 0.061	0.156	1.030 ± 0.134		
6	0.26	0.240 ± 0.018	0.5	-0.092 ± 0.053	0.130	1.065 ± 0.115		
7	0.26	0.226 ± 0.012	0.6	-0.104 ± 0.073	0.104	1.082 ± 0.142		
8	0.26	0.225 ± 0.022	0.7	-0.084 ± 0.150	0.078	1.043 ± 0.208		
9	0.26	0.207 ± 0.029	0.8	0.004 ± 0.453	0.052	0.989 ± 0.370		
10	0.26	0.143 ± 0.028	0.9	-0.151 ± 0.043	0.026	1.097 ± 0.068		

TABLE III. Estimated parameters n_e , A^* , and K^* with the SDM compared with the true parameters n = 0.26, α , and K of the ETAS model used to generate the synthetic catalogs by the generation-by-generation algorithm. The other parameters have been fixed to $m_d = m_0 = 0$, b = 1, c = 0.002, and $\theta = 0.19$.

declustered catalogs (pure Poisson) misclassifies about 23% of the events as being triggered, while they are all exogenous.

5. Tests using catalogs generated by the generation-bygeneration algorithm B

Test of the SDM. We tested the SDM using data sets generated with various values of α =0–0.9 of productivity law (9). We generated ten catalogs of length of ~50 000 days for each parameters set, with a background rate μ_0 =1 per day. The duration of the catalogs in terms of days is just to offer a convenient interpretation, as the intrinsic time scale is more generally determined by 1/ μ_0 . Each catalog contained about 70 000–100 000 events. We removed the first 2000 events in the first part of the catalogs, roughly corresponding to the first 1500 days, and applied the SDM procedure 20 times to each catalog. Table III compiles the obtained values n_e , A^* , and K^* defined in Eqs. (19) and (20), reports their standard deviations, and compares with the true values $n_{,,} \alpha$, and K.

We found that the distributions of background rates for the different values of $\alpha = 0-0.9$ are reasonably estimated, with an error of no more than 10%. While the estimated branching ratio n_e is also reasonably close to the true value for α up to 0.6 and then starts to systematically deviate for larger α 's, the estimated parameters K^* and A^* are found very far from the true values. In particular, the values of the estimated fertility exponent A^* would imply that events of all magnitudes have on average no more than one aftershock, which is very far from being the case, especially for large values of α . One partial cause for this bad result is the Omori law which, as shown in Fig. 1, implies very long interevent times between direct aftershocks. Some of those intervals can be longer than our catalog and the real productivity will be thus underestimated. Another possible partial cause is that the removal of an initial part of the catalogs to analyze the more stationary regime at later times deletes by construction many events that are mothers of the observed events. This also leads to an underestimation of the triggering productivity. These two explanations suggest that the problem is intrinsic to the application of the SDM to the ETAS model and we do not envision easy fixes.

Test of the MISD method. We applied the MISD method to the same catalogs. Table IV shows a slight systematic overestimation of the branching ratio n_e over the true value n.

TABLE IV. Estimated parameters n_e , A^* , and K^* with the MISD compared with the true parameters n = 0.26, α , and K of the ETAS model used to generate the synthetic catalogs by the generation-by-generation algorithm. The other parameters have been fixed to $m_d = m_0 = 0$, b = 1, c = 0.002, and $\theta = 0.19$.

No.	n	n _e	α	A^*	K	K^*
1	0.26	0.319 ± 0.048	0.0	0.332 ± 0.265	0.260	0.931 ± 0.450
2	0.26	0.299 ± 0.035	0.1	0.465 ± 0.535	0.234	1.077 ± 1.106
3	0.26	0.328 ± 0.064	0.2	0.352 ± 0.264	0.208	0.925 ± 0.693
4	0.26	0.343 ± 0.061	0.3	0.488 ± 0.373	0.182	0.847 ± 1.047
5	0.26	0.336 ± 0.062	0.4	0.355 ± 0.134	0.156	0.955 ± 0.678
6	0.26	0.332 ± 0.069	0.5	0.579 ± 0.265	0.130	0.373 ± 0.372
7	0.26	0.322 ± 0.060	0.6	1.011 ± 0.912	0.104	0.293 ± 0.325
8	0.26	0.320 ± 0.113	0.7	0.815 ± 0.488	0.078	0.295 ± 0.312
9	0.26	0.337 ± 0.151	0.8	0.818 ± 0.223	0.052	0.122 ± 0.128
10	0.26	0.320 ± 0.174	0.9	0.900 ± 0.449	0.026	0.126 ± 0.143

The estimations of α and *K* are also bad with large systematic errors. The trend of variation in *K* as a function of α for the fixed true n=0.26 is qualitatively reproduced by the dependence of K^* as a function of α .

We must also report a surprising difference between the SDM and the MISD method. While the standard deviations of the estimated parameters over ten different synthetic catalogs are sometimes significantly larger for the MISD method compared with the SDM, the former method exhibited sometimes very accurate results for a few catalogs for some specific values of the parameters. For instance, for one of the synthetic catalog generated with the parameters n=0.26, α =0.9, K=0.026, and μ_0 =1, the MISD method gave the following estimates: $n_e = 0.255 \pm 0.007$, $A^* = 0.916 \pm 0.037$, K^* $=0.018 \pm 0.007$, and background rate $\hat{\mu}_0 = 0.890 \pm 0.218$. The existence of such an excellent inversion has to be tempered by the fact that the estimates obtained with the MISD method applied to the other nine catalogs generated with the same parameters were bad. This suggests a very strong dependence of the performance of the MISD method on the specific stochastic realizations.

Impact of catalog incompleteness. We now report some results on the estimation of parameters by the two declustering methods applied to incomplete catalogs, motivated by the nature of real-life catalogs. The incompleteness is measured by the magnitude threshold $m_d > m_0$, below which events are missing from the catalogs used for the SDM and MISD method.

First, we focus on the results [42,43] that the branching ratio n is renormalized into an effective value n_t which is a decreasing function of m_d ,

$$n_t(m_d) = \frac{1}{1 + \frac{1 - n}{n} [10^{\alpha m_d}]^{b/\alpha - 1}}.$$
 (22)

This prediction was verified by direct simulations with the ETAS model. Here, we test how the incompleteness of catalogs may interfere with the stochastic declustering methods. We generated 50 catalogs with the following set of parameters: b=1, c=0.001, $\theta=0.1$, n=0.7, $\alpha=0.7$, and K=0.21 and varied m_d from 0 to 2. Applying the SDM 20 times to each of the 50 synthetic data sets, Fig. 2 shows that (i) the general trend of decreasing $n_t(m_d)$ as a function of m_d is recovered, but (ii) there is a very significant downward bias of approxi-



FIG. 2. Dependence of the effective branching ratio n_t as a function of the threshold magnitude m_d of catalog incompleteness. The parameters used to generate the synthetic catalogs with the ETAS model are b=1, c=0.001, n=0.7, $\alpha=0.7$, and K=0.21. The dashed (respectively dotted) line corresponds to n_t obtained by using SDM algorithm (MISD algorithm). The continuous curve is the validated theoretical formula (22).

mately 0.2 over the whole range $0 \le m_d \le 2$. Table V shows unsurprisingly that the estimated parameters A^* and K^* are very far from the true values α and K. Using incomplete catalogs cannot be expected to improve the estimation of parameters which is already bad for complete catalogs. For smaller values of the true branching ratio n, the discrepancy is smaller between the reconstructed n_t and the theoretical formula (22). For instance, for n=0.4, the difference between the estimated n_t and the formula (22) decreases from 0.1 for $m_d=0$ (no incompleteness) to almost zero for $m_d=2$ for which $n_t(m_d=2) \approx 0.04$.

Typical results for the MISD method are reported in Table VI for the set of true parameters n=0.7, $\alpha=0.7$, and K=0.21, for m_d varying from 0 to 2. While the estimated n_e for $m_d=0$ (no incompleteness) are as good as with the direct SDM estimation method, the other estimated parameters are strongly biased. For incomplete catalogs $m_d>0$, we observe a large overestimation of n_t and significant errors in the other estimated parameters.

Another test is provided by comparing the CDFs of background events in the incomplete catalogs as a function of m_d obtained by the declustering methods with the true CDFs. We found that the larger is the true branching ratio n, the

TABLE V. Estimated parameters n_e , A^* , and K^* with the SDM compared with the true parameters n = 0.7, $\alpha = 0.7$, and K = 0.21 of the ETAS model used to generate the synthetic catalogs by the generation-by-generation algorithm, for various magnitude threshold m_d of incompleteness. The other parameters have been fixed to b=1, c=0.001, and $\theta=0.1$. N_{m_d} is the number of events in the incomplete catalogs.

No.	m_d	N_{m_d}	n _e	A^*	К*
1	0.0	2000	0.488 ± 0.032	-0.045 ± 0.400	1.037 ± 0.214
2	0.5	2000	0.406 ± 0.035	0.135 ± 0.784	1.085 ± 0.437
3	1.0	1000	0.319 ± 0.034	-0.006 ± 0.286	1.201 ± 0.491
4	1.5	500	0.271 ± 0.056	-0.054 ± 0.116	1.376 ± 0.574
5	2.0	150	0.204 ± 0.067	-0.047 ± 0.094	1.495 ± 0.769

TABLE VI. Estimated parameters n_e , A^* , and K^* with the MISD method compared with the true parameters n=0.7, $\alpha=0.7$, and K=0.21 of the ETAS model used to generate the synthetic catalogs by the generation-by-generation algorithm, for various magnitude threshold m_d of incompleteness. The other parameters have been fixed to b=1, c=0.001, and $\theta=0.1$. N_{m_d} is the number of events in the incomplete catalogs.

No.	m_d	N_{m_d}	n _e	A^*	<i>K</i> *
1	0.0	4000	0.486 ± 0.074	0.930 ± 0.448	0.167 ± 0.166
2	0.5	1300	0.420 ± 0.069	0.873 ± 0.878	0.089 ± 0.188
3	1.0	400	0.345 ± 0.060	-0.357 ± 2.804	0.239 ± 0.497
4	1.5	130	0.327 ± 0.089	0.249 ± 1.521	0.168 ± 0.686
5	2.0	40	0.296 ± 0.097	-0.771 ± 3.214	0.002 ± 0.874

larger is the discrepancy between the true and reconstructed CDFs of background events, using both declustering methods for all m_d values. For $n \le 0.4$, the reconstructed background CDFs are in reasonable agreement with the theoretical formula (22) with typical errors of about 10% (see Tables IX and V and Fig. 3). For large branching ratios, the errors are too large and the declustering methods are unreliable.

Comparing Figs. 2 and 3, one can notice that MISD becomes less precise for large thresholds m_d . That is more likely caused by the smaller lengths of the catalogs.

6. Tests using catalogs generated by the event-by-event algorithm IIIC

The same tests as reported in the previous subsection were performed on catalogs generated by the event-by-event algorithm IIIC. We recall that our motivation for using this alternative algorithm is to test for the expected influence of transient regimes, which are absent by construction in the synthetic catalogs obtained with the event-by-event algorithm IIIC.

Using one hundred catalogs of 2000 events (ten for each α going from 0 to 0.9) generated with algorithm IIIC, we



FIG. 3. Dependence of the effective branching ratio n_t as a function of the threshold magnitude m_d of catalog incompleteness. The parameters used to generate the synthetic catalogs with the ETAS model are b=1, c=0.001, $\theta=0.1$, n=0.4, $\alpha=0.2$, K=0.32. The dashed (dotted) line corresponds to n_t obtained by using the SDM algorithm (MISD algorithm). The continuous curve is the validated theoretical formula (22).

applied the SDM and obtained the results summarized in Table VII. The results are similar to those of Table III, with a reasonable estimation of n_e but estimated A^* and K^* are very far from the true values α and K.

Ten catalogs of 4000 events were generated with the parameters b=1, c=0.001, $\theta=0.1$, n=0.6, $\alpha=0.2$, and K=0.48. Incompleteness was introduced at magnitude thresholds m_d varying from 0 to 2, reducing the size of catalogs to the m_d -dependent number $\approx N_{m_d}$ events. Applying the SDM to these incomplete catalogs, we obtained the results shown in Table VIII. As mentioned above, the estimated n_e is found in better agreement with theoretical prediction (22) for the largest m_d values for which n_t is the smallest.

Table IX shows that the MISD method gives results similar to those previously obtained in Table VI. The estimated value of background rate is $\hat{\mu}_0 = 1.11 \pm 0.15$, which is close to the true value 1.

V. TESTS ON THE INFLUENCE OF MEMORY (EXPONENT θ), BACKGROUND RATE AND CATALOG LENGTHS

A. Influence of the value of the memory exponent θ

We now test one possible origin for the rather bad performance of the SDM and MISD method when using the thinning probabilities, namely, the very long memory quantified by the small value of the exponent θ [as defined in expression (8)] used in the simulations.

A series of tests were made with a larger value of $\theta = 0.5$, corresponding to significantly shorter memory. We varied the values of *n* and α , while fixing the other parameters *c* =0.001, b=1, and $\mu=1$. We generated ten catalogs of 2500 events each using the event-by-event algorithm described in Sec. III C for each set of parameters. We implemented both declustering algorithms 20 times to each catalog. Table X presents the resulting estimates of the parameters. The most striking result is that the branching ratio *n* estimated by the SDM is in general very good. While the MISD method is not reliable for estimating the branching ratio n, it is better for the estimation of the productivity exponent α , especially for large values. The results presented in Table X, when put in comparison with those of the previous tables obtained with much smaller values of the memory exponent θ , illustrate clearly the impact of the long memory on the declustering results.

argonum. The other parameters are fixed to $m_d = 0, b = 1, c = 0.002$, and $b = 0.17$.								
No.	п	n _e	α	A^*	K	Κ*		
1	0.26	0.244 ± 0.024	0.0	0.042 ± 0.092	0.260	0.845 ± 0.211		
2	0.26	0.238 ± 0.027	0.1	0.033 ± 0.082	0.234	0.856 ± 0.203		
3	0.26	0.248 ± 0.022	0.2	0.016 ± 0.077	0.208	0.895 ± 0.176		
4	0.26	0.246 ± 0.036	0.3	-0.014 ± 0.075	0.182	0.940 ± 0.173		
5	0.26	0.233 ± 0.011	0.4	-0.032 ± 0.070	0.156	0.967 ± 0.137		
6	0.26	0.234 ± 0.024	0.5	-0.025 ± 0.072	0.130	0.960 ± 0.152		
7	0.26	0.233 ± 0.017	0.6	-0.101 ± 0.065	0.052	1.069 ± 0.107		
8	0.26	0.216 ± 0.024	0.7	-0.074 ± 0.060	0.104	1.011 ± 0.116		
9	0.26	0.206 ± 0.021	0.8	0.179 ± 0.952	0.078	1.005 ± 0.314		
10	0.26	0.216 ± 0.210	0.9	-0.121 ± 0.055	0.026	1.151 ± 0.205		

TABLE VII. Estimated parameters n_e , A^* , and K^* with the SDM compared with the true parameters n = 0.26, α , and K of the ETAS model used to generate the synthetic catalogs using the event-by-event algorithm. The other parameters are fixed to $m_d=0$, b=1, c=0.002, and $\theta=0.19$.

B. Case of a single large main shock triggering aftershocks

In another series of tests, we check if the SDM or MISD are able to determine a pure tree branching process emanating from a single source. In this goal, we removed all spontaneous events except one, the first "main shock." This corresponds to imposing $\mu=0$. In order to have sufficiently many events in the catalogs, we take the magnitude of the single main shock large $(M_1=7)$ and also impose large values for the parameters *n* and α to produce long sequences of events. The parameter θ was taken equal to 0.1 (long memory). We found that the SDM recognized the existence of the only background event in two tests out of three (for $\alpha = 0.8$). In contrast, the MISD lacks efficiency in such conditions and proposes significantly nonzero values for the background rate μ . The estimation of *n* with formula (19) cannot provide good results because it should give $n_e \rightarrow 1$ (only one background event) for large catalogs. Other methods of estimating the branching ratio described in [12] gave even worse results, with $n_e > 1$.

C. Influence of catalog lengths

As mentioned above, two effects combine to limit the efficiency of the SDM and MISD methods: the smallness of the exponent θ leading to very long memory and the limited length of the catalogs. We investigated how these two effects

are inter-related in practice. Using algorithm IIIB, we generated 100 catalogs with variable numbers of events (from 3000 to 12 000) and for various values of θ (from 0.05 up to 0.5). The other parameters were fixed at b=1, $\alpha=0.7$, c=0.001, and n=0.7. We then applied the mSDM to each catalog. Figure 4 shows the estimation error of the branching ratio *n*. One can observe very large variations from realization to realization and a weak tendency for estimation errors to decrease with the length of the catalogs.

To be more quantitative, let us introduce the cumulative error ratio defined as the sum of squares of relative errors over all parameters,

$$\varepsilon = \sum_{i=1}^{5} \left(1 - \frac{\Theta_i^d}{\Theta_i^s} \right)^2, \quad (\Theta_{1-5} = b, \alpha, c, \theta, n).$$
(23)

The dependence of the error in the determination of the exponent θ as a function of catalog lengths is shown in Fig. 5 for a number of realizations for two classes of catalogs sorted according to their overall estimation error measured by ε as defined by Eq. (23). Again, one can observe an overall decrease in the estimation error with the length of the catalogs, decorated by a very large variability from catalog to catalog.

As an illustration of the quality of the reconstruction of the tree structure of ancestry in different catalogs, we took seven catalogs with $\varepsilon < 0.1$ (best results) and several catalogs with $\varepsilon > 1$ (bad results) and determined the percentage of

TABLE VIII. Estimated parameters n_e , A^* , and K^* with the SDM compared with the true parameters n = 0.6, $\alpha = 0.2$ and K = 0.48 of the ETAS model used to generate the synthetic catalogs by the event-by-event algorithm, for various magnitude threshold m_d of incompleteness. The other parameters have been fixed to b=1, c=0.001, and $\theta=0.1$.

No.	m_d	N_{m_d}	n _e	A^*	<i>K</i> *
1	0.0	4000	0.445 ± 0.012	-0.004 ± 0.048	0.960 ± 0.115
2	0.5	1300	0.208 ± 0.012	0.000 ± 0.0345	0.968 ± 0.098
3	1.0	400	0.080 ± 0.014	-0.004 ± 0.022	1.013 ± 0.085
4	1.5	130	0.024 ± 0.012	0.003 ± 0.017	0.977 ± 0.073
5	2.0	40	0.003 ± 0.009	0.009 ± 0.018	0.932 ± 0.077

TABLE IX. Estimated parameters n_e , A^* , and K^* with the MISD method compared with the true parameters n=0.4, $\alpha=0.2$, and K=0.32 of the ETAS model used to generate the synthetic catalogs by the generation-by-generation algorithm, for various magnitude threshold m_d of incompleteness. The other parameters have been fixed to b=1, c=0.001, and $\theta=0.1$.

No.	m _d	N_{m_d}	n _e	A^*	<i>K</i> *
1	0.0	2000	0.365 ± 0.082	0.380 ± 0.198	0.618 ± 0.393
2	0.5	2000	0.235 ± 0.079	0.275 ± 1.462	0.324 ± 0.450
3	1.0	1000	0.170 ± 0.077	0.041 ± 2.094	0.348 ± 0.597
4	1.5	500	0.176 ± 0.074	0.065 ± 1.501	0.364 ± 0.651
5	2.0	150	0.193 ± 0.126	0.114 ± 0.991	0.411 ± 0.621

background events recognized as background and of aftershocks recognized as aftershocks. For catalogs with the best directly estimated parameters, the percentage were 66% and 72%, respectively, while for the "bad" catalogs the results were 9% and 93% (i.e., almost all events were incorrectly recognized as aftershocks).

VI. CONCLUSIONS

Many time series in natural and social sciences can be seen as embodying an interplay between exogenous influences and an endogenous organization. We have used a simple model of events occurring sequentially, in which future events are influenced (partially triggered) by past events to ask the question of how well can one disentangle the exogenous events from the endogenous ones.

The exogenous events are modeled here by a Poisson flow of so-called background events with constant intensity μ_0 . The ETAS specification of the conditional self-excited Hawkes Poisson model has been used. It contains three principal ingredients: (i) a long Omori-like power-law memory of the influence of past events on future events, (ii) a Gutenberg-Richter-like distribution of event magnitudes, and (iii) a fertility law expressing how many events are triggered by a given event as a function of its magnitude.

In order to separate background events from triggered events, we have implemented and compared two so-called "declustering" algorithms, the SDM introduced by Zhuang *et al.* [22,23], and the MISD method proposed by Marsan and Lengliné [25]. We have applied these two methods to synthetic catalogs generated by two algorithms using the ETAS model, the generation-by-generation algorithm IIIB, and the event-by-event algorithm IIIC.

We specifically address the problem of reconstructing the tree of ancestry of the sequence of events in recorded catalogs. In particular, one main goal is to distinguish the exogenous shocks (background events) from the endogenous (triggered events). For these problems, we find that declustered catalogs obtained from synthetic catalogs generated with the ETAS model are rather unreliable, when using catalogs with of the order of thousands of events, typical of realistic applications. The estimated rates of exogenous events suffer from large errors. The key branching ratio n, quantifying the fraction of events that have been triggered by previous events, is also badly estimated in general with these approaches. We find however that the errors tend to be smaller and perhaps acceptable in some cases for the smaller fertility exponent $\alpha < 0.6$ and for the smaller branching ratio n < 0.4 typically. Results become better when the memory exponent θ is increased, i.e., when the memory is shortened. We do not find significantly better performance of the SDM versus the MISD method or vice versa, with the curious observation that the MISD method is sometimes very precise for some catalog realizations, but this property is not robust

TABLE X. Estimated parameters n_e , A^* , and K^* with the SDM and MISD compared with the true parameters n, α , and K of the ETAS model used to generate the synthetic catalogs by the event-by-event algorithm. The other parameters have been fixed to $m_d=m_0=0$, b=1, c=0.001, and $\theta=0.5$.

No.	Method	n	α	K	n _e	A^*	K^*
1	SDM	0.2	0.2	0.16	0.199 ± 0.009	-0.020 ± 0.042	0.999 ± 0.087
2	MISD	0.2	0.2	0.16	0.313 ± 0.019	1.116 ± 1.203	0.367 ± 0.285
3	SDM	0.5	0.5	0.25	0.502 ± 0.020	-0.061 ± 0.076	1.004 ± 0.156
4	MISD	0.5	0.5	0.25	0.546 ± 0.018	0.970 ± 0.828	0.501 ± 0.363
5	SDM	0.8	0.8	0.16	0.698 ± 0.072	0.024 ± 0.481	0.941 ± 0.250
6	MISD	0.8	0.8	0.16	0.755 ± 0.045	0.700 ± 0.108	0.347 ± 0.209
7	SDM	0.8	0.2	0.64	0.793 ± 0.014	0.009 ± 0.068	0.942 ± 0.141
8	MISD	0.8	0.2	0.64	0.804 ± 0.016	0.352 ± 0.337	0.626 ± 0.272
9	SDM	0.2	0.8	0.04	0.168 ± 0.021	-0.160 ± 0.081	1.187 ± 0.159
10	MISD	0.2	0.8	0.04	0.259 ± 0.037	0.694 ± 0.128	0.175 ± 0.132



FIG. 4. Error of the direct estimation of the branching ratio by the mSDM as a function of catalog length for different values of θ .

with respect to other stochastic realizations. We have also investigated the role of incompleteness on declustering, and found that this is not the essential limiting problem.

We should however make clear that these rather negative results are not necessarily opposed to the more positive results reported by Zhuang *et al.* and Marsan and Lengliné, which refer to a different problem, that of the direct estimation of the model parameters (and not of the tree of ancestry and of the distinction between exogenous and endogenous events). Our larger ambition to reconstruct the tree of ancestry has identified clearly intrinsic limits of the inversion process.

It appears that the high level of randomness together with the very long memory makes the stochastic reconstruction of trees of ancestry and the estimation of the key parameters quite unreliable. Technically, we find the coexistence of many coexisting stochastic reconstructions with different parameter estimates, and it is not obvious how to select what



FIG. 5. Error in the determination of the exponent θ as a function of catalog lengths is shown for two classes of catalogs sorted according to their overall estimation error measured by ε as defined by Eq. (23). The straight line is helping the eye defining a lower envelop of the cloud of points.

should be the right one. This question is reminiscent of complex optimization problem in the presence of a very large number of almost equivalent solutions, as occurs in so-called NP-complete problems [45]. There thus appears to be fundamental limitations intrinsic to this class of models.

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- P. Brémaud, *Point Processes and Queues* (Springer, New York, 1981).
- [2] D. J. Daley and D. Vere-Jones, An Introduction to the Theory of Point Processes: Elementary Theory and Methods, 2nd ed., (Springer, New York, 2003), Vol. 1.
- [3] D. Ruelle, Phys. Today 57(5), 48 (2004).
- [4] D. Sornette, in *Extreme Events in Nature and Society*, edited by S. Albeverio, V. Jentsch, and H. Kantz (Springer, Berlin, 2005), pp. 95–119.
- [5] D. Sornette, F. Deschatres, T. Gilbert, and Y. Ageon, Phys. Rev. Lett. 93, 228701 (2004).
- [6] F. Deschatres and D. Sornette, Phys. Rev. E **72**, 016112 (2005).
- [7] B. M. D. Roehner, Int. J. Mod. Phys. C 15, 809 (2004).
- [8] D. Sornette, Y. Malevergne, and J. Muzy, Risk 16, 67 (2003).
- [9] D. Sornette, Why Stock Markets Crash (Critical Events in Complex Financial Systems) (Princeton University Press, New

Jersey, 2003).

- [10] A. Johansen and D. Sornette, Brussels Economic Review (Cahiers economiques de Bruxelles) 49(3/4) (2006); e-print arXiv:cond-mat/0210509.
- [11] A. Helmstetter, D. Sornette, and J.-R. Grasso, J. Geophys. Res. 108, 2046 (2003).
- [12] A. Helmstetter and D. Sornette, Geophys. Res. Lett. **30**, 2069 (2003).
- [13] D. Sornette, V. I. Yukalov, E. P. Yukalova, J.-Y. Henry, D. Schwab, and J. P. Cobb, J. Biol. Syst. 17, 225 (2009).
- [14] I. Osorio, M. G. Frei, D. Sornette, J. Milton, and Y.-C. Lai, e-print arXiv:0712.3929.
- [15] Y. Ogata, J. Am. Stat. Assoc. 83, 9 (1988).
- [16] A. Helmstetter and D. Sornette, J. Geophys. Res. 107, 2237 (2002).
- [17] A. Helmstetter and D. Sornette, J. Geophys. Res., [Solid Earth] 108, 2457 (2003).

- [18] A. G. Hawkes, J. R. Stat. Soc. Ser. B (Methodol.) 33, 438 (1971).
- [19] A. G. Hawkes and L. Adamapoulos, Bull. Internat. Statist. Inst. 45, 454 (1973).
- [20] A. G. Hawkes and D. Oakes, J. Appl. Probab. 11, 493 (1974).
- [21] A. Saichev and D. Sornette, J. Geophys. Res. 112, B04313 (2007).
- [22] J. Zhuang, Y. Ogata, and D. Vere-Jones, J. Am. Stat. Assoc. 97, 369 (2002).
- [23] J. Zhuang, Ph.D. thesis, Department of Statistical Science, The Graduate University for Advanced Studies, 2002.
- [24] J. Zhuang, Y. Ogata, and D. Vere-Jones, J. Geophys. Res. 109, B05301 (2004).
- [25] D. Marsan and O. Lengliné, Science 319, 1076 (2008).
- [26] Y. Ogata, IEEE Trans. Inf. Theory IT-27, 23 (1981).
- [27] MISD (http://www.lgit.univ-savoie.fr/MISD/misd.html).
- [28] P. Brémaud and L. Massoulié, Ann. Probab. 24, 1563 (1996).
- [29] P. Brémaud and L. Massoulié, J. Appl. Probab. 38, 122 (2001).
- [30] D. Sornette and G. Ouillon, Phys. Rev. Lett. **94**, 038501 (2005).
- [31] G. Ouillon and D. Sornette, J. Geophys. Res. 110, B04306 (2005).
- [32] A. Saichev and D. Sornette, Phys. Rev. E 74, 011111 (2006).
- [33] D. Sornette, S. Utkin, and A. Saichev, Phys. Rev. E 77,

066109 (2008).

- [34] K. R. Felzer, T. W. Becker, R. E. Abercrombie, G. Ekström, and J. R. Rice, J. Geophys. Res. 107, 2190 (2002).
- [35] T. Ozaki, Ann. Inst. Stat. Math. 31, 145 (1979).
- [36] A. Saichev and D. Sornette, Phys. Rev. Lett. 97, 078501 (2006).
- [37] J.-P. Bouchaud and A. Georges, Phys. Rep. 195, 127 (1990).
- [38] J. Beran, Statistics for Long-Memory Processes (Monographs on Statistics and Applied Probability) (Chapman and Hall, London, 1994).
- [39] R. Metzler and J. Klafter, Phys. Rep. 339, 1 (2000).
- [40] D. Sornette, Critical Phenomena in Natural Sciences, 2nd ed., Springer Series in Synergetics (Springer, Heidelberg, 2006).
- [41] Y. Ogata and J. Zhuang, Tectonophysics 413, 13 (2006).
- [42] D. Sornette and M. J. Werner, J. Geophys. Res. 110, B09303 (2005).
- [43] A. Saichev and D. Sornette, Eur. Phys. J. B 51, 443 (2006).
- [44] J. Zhuang, A. Christophersen, M. K. Savage, D. Vere-Jones, Y. Ogata, and D. D. Jackson, J. Geophys. Res. 113, B11302 (2008).
- [45] M. R. Garey and D. S. Johnson, Computers and Intractability: A Guide to the Theory of NP-Completeness (Freeman, New York, 1979).